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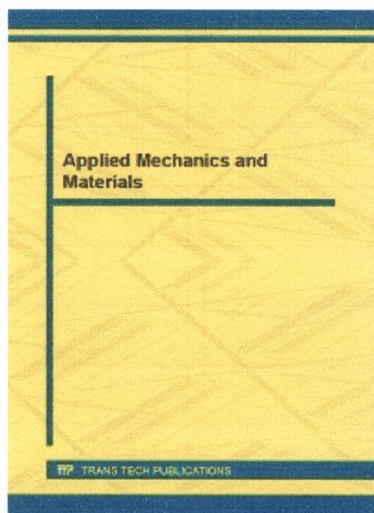
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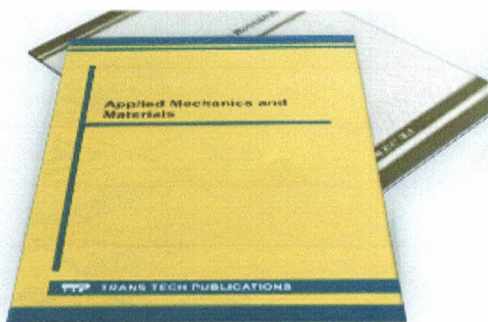
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In this volume collected articles presented at Uttaradit Rajabhat University International Conference on Science and Technology 2018 (URUICST 2018. August 2-3, 2018, Uttaradit Thailand). The main topics of issue are technologies in the food processing and agriculture, pharmacology and biotechnologies, technologies for production of renewable energy, advanced materials and chemical technologies in the environmental engineering, designing in mechanical engineering and information technologies.

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Determine the Specific Heat Capacity, the Entropy, and the Standard Deviation of Particle Vibration Bound-State under the Anharmonic Oscillator Asymmetric Potential by Partition Function Method

Metaporn Apiratigosol^{1,a}, Katang Jeeradit^{1,b}, Pattadon Keawpeai^{1,c}
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Keywords: partition function, specific heat capacity, entropy, standard-deviation, anharmonics oscillator asymmetric potential

Abstract. In this work, we consider is the single-particle bound-state problem. A single-particle moves in the absence of the anharmonics oscillator asymmetric potential. The bound-state problem then is to solve the partition function for single-particle in the presence of anharmonics oscillator asymmetric potential. The factor second exponential function of partition function of particle bound in the anharmonics oscillator asymmetric potential can also be expanded in power series of the temperature and parameter b , c , d and x is the position of single-particle bound-state problem and using the integrate Gaussian. The purpose of this works, we will compute the partition function as a function of temperature, entropy, specific heat capacity and standard deviation which depend on the c, d parameter and frequency (f).

Introduction

In general, the macroscopic world consists of many particles simultaneously interacting with one another. Currently, it had shown the direct interaction algorithm (DIA) for approximating the canonical partition function Z_N in N^2 operations and the DIA was used to compute equilibrium properties for the isotropic 2D Ising model, and the accuracy of the DIA was compared to that of basic Metropolis Monte Carlo method [3]. And the numerically studied the standard canonical sampling of phase space to directly compute the partition function and heat capacity [4]. In addition a method for calculating partition function taking into account anharmonic contributions for system involving both small-amplitude vibration and hindered rotations was present by Skouteris et. al.[6]. The purpose of this works, we will compute the partition function as a function of temperature, internal energy, entropy, specific heat capacity, enthalpy, free energy and standard deviation. The scheme of the article is as follows. In section (2) we write the basic evaluation of integral Gaussian odd and even function. In section (3) we illustrate evaluation of the partition function, internal energy, entropy, specific heat capacity, enthalpy, free energy and standard deviation of particle bound state in anharmonics oscillator asymmetric potential. In section (4) contains our conclusions.

Evaluation of Integral

An important example of this method is provided by the evaluation of integral

$$I = \int_{-\infty}^{\infty} e^{-x^2} dx \quad (1)$$

Its value may be found by first constructing I^2 , as follow :

$$I^2 = \int_{-\infty}^{\infty} e^{-x^2} dx \int_{-\infty}^{\infty} e^{-y^2} dy = \iint_R e^{-(x^2+y^2)} dx dy, \quad (2)$$

where the region R is the whole xy -plane. The differential $dx dy$ represents an element of area in cartesian coordinates, with the domain of integration extending over the entire xy -plane. An alternative representation of the last integral can be expressed in plane polar coordinates r, θ . The two coordinate systems are related by $x = r \cos \theta$, $y = r \sin \theta$. The element of area in polar coordinates is given by $r dr d\theta$, so that the double integral becomes

$$I^2 = \iint_R e^{-\rho^2} \rho d\rho d\phi = \int_0^{2\pi} d\phi \int_0^{\infty} \rho e^{-\rho^2} d\rho = 2\pi \left[-\frac{1}{2} e^{-\rho^2} \right]_0^{\infty} = \pi.$$

Therefore the original integral is given by $I = \sqrt{\pi}$. We note, however, that unlike in all the previous examples, the regions of integration R' and R'' are both infinite in extent. It is therefore prudent to derive this result more rigorously; this we do by considering the integral

$$\int_{-\infty}^{\infty} e^{-ax^2} dx = \frac{1}{\sqrt{a}} \int_{-\infty}^{\infty} e^{-ay^2} dy = \sqrt{\frac{\pi}{a}}, \quad (3)$$

and integrate gaussian even function

$$\int_{-\infty}^{\infty} x^2 e^{-ax^2} dx = -\int_{-\infty}^{\infty} \frac{\partial}{\partial a} e^{-ax^2} dx = -\frac{d}{da} \int_{-\infty}^{\infty} e^{-ax^2} dx = -\frac{d}{da} \sqrt{\frac{\pi}{a}} = \frac{1}{2} \sqrt{\pi a^{-\frac{3}{2}}}, \quad (4)$$

and integrate gaussian odd function

$$\int_{-\infty}^{\infty} x^2 e^{-ax^2} dx = 0 \quad (5)$$

and integrate gaussian even function

$$\int_{-\infty}^{\infty} x^4 e^{-ax^2} dx = -\int_{-\infty}^{\infty} \frac{\partial^2}{\partial a^2} e^{-ax^2} dx = -\frac{d^2}{da^2} \int_{-\infty}^{\infty} e^{-ax^2} dx = \frac{4}{4} \sqrt{\pi a^{-\frac{5}{2}}}, \quad (6)$$

and

$$\int_{-\infty}^{\infty} x^6 e^{-ax^2} dx = -\int_{-\infty}^{\infty} \frac{\partial^3}{\partial a^3} e^{-ax^2} dx = \frac{\partial^2}{\partial a^2} \left(-\int_{-\infty}^{\infty} x^2 e^{-ax^2} dx \right) = \frac{15}{8} \sqrt{\pi a^{-\frac{7}{2}}}, \quad (7)$$

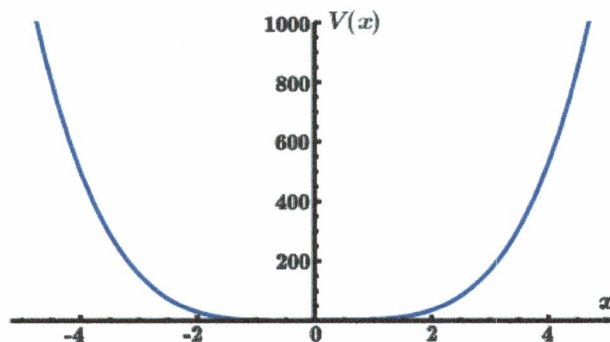


Fig. 1: The anharmonics oscillator asymmetric potential system.

Let us consider a problem to illustrate the use of integrate Gaussian function(see eq4-eq7). An interesting work is that of a particle move under the influence of the anharmonics oscillator asymmetric potential in section (3).

Evaluation of the Partition Function of Particle Bound State in Anharmonics Oscillator Asymmetric Potential

In this section, we consider the statistical mechanics behavior of a particle that moves under the influence of the anharmonics oscillator asymmetric potential as

$$V(x) = \frac{1}{2}m\omega^2 x^2 + bx + cx^3 + dx^4 \quad (8)$$

where b , c , and d is parameter positive constant. The problem then is to solve the partition function for a particle in the presence of the anharmonics oscillator asymmetric potential see figure (1). The partition function of particle bound in the anharmonics oscillator asymmetric potential is

$$Z(T) = \int_{-\infty}^{\infty} e^{-\frac{m\omega^2}{2kT}x^2} e^{-\frac{1}{kT}(bx + cx^3 + dx^4)} dx \quad (9)$$

Formally, if we expand the exponential function in a Taylor series about the minimum:

$$e^{-\frac{1}{kT}(bx + cx^3 + dx^4)} = 1 - \frac{bx}{kT} + \frac{b^2 x^2}{2(kT)^2} - \frac{cx^3}{kT} + \left(\frac{bc}{(kT)^2} - \frac{d}{kT} \right) x^4 + \frac{bd}{(kT)^2} x^5 + \frac{c^2}{2(kT)^2} x^6 + \dots \quad (10)$$

Since the coefficients b , c , and d are small, the integral is largely determined by the first exponential function, which is appreciable only in the range $x^2 \leq 2kT / m\omega^2$. In this range,

$$cx^3 / kT \leq c\sqrt{kT} / (m\omega^2 / 2)^{\frac{1}{2}} \ll 1 \text{ and } dx^4 / kT \leq dkT / (m\omega^2 / 2)^2 \ll 1.$$

Substituting equation (10) and using integrate gaussian of equation (4)-(7) into equation (9), we can rewrite the partition function completely in terms of temperature as

$$Z(T) = \sqrt{\frac{2\pi kT}{m\omega^2}} \left\{ 1 + \frac{b^2}{2mkT\omega^2} + \frac{3bc}{m^2\omega^4} - \frac{12d}{4m^2\omega^4}(kT) + \frac{15c^2}{2m^3\omega^6}(kT) + \frac{105d^2}{2m^4\omega^8}(kT)^2 + \dots \right\} \quad (11)$$

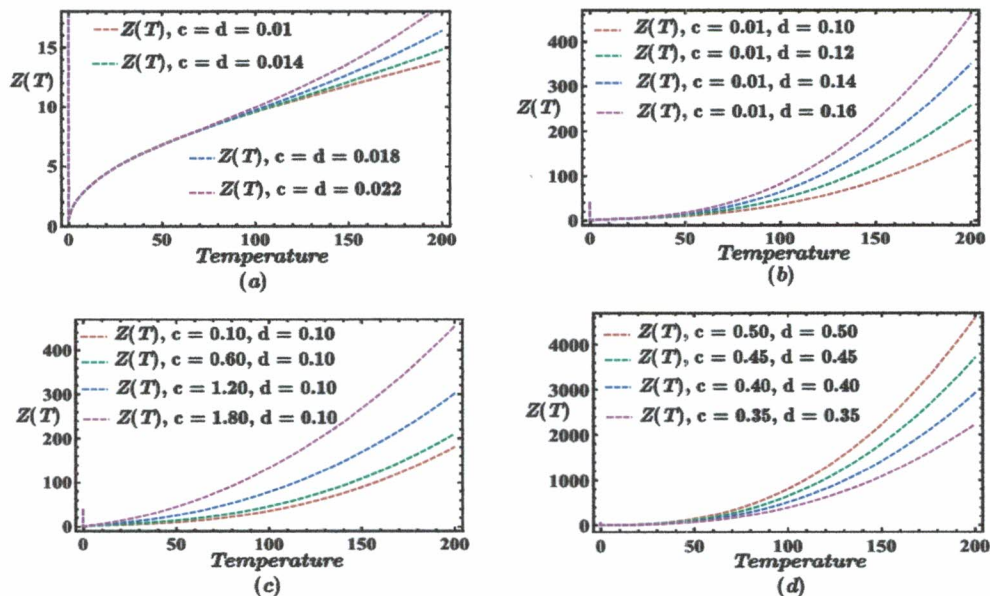


Fig. 2: Partition function as a function of temperature plotted for c and d is parameter. positive constant.

Schematically drawn, the partition function in equation (11) has the shape shown in figure (2). Hence from knowledge of the partition function $Z(T)$, the internal energy per oscillation $U(T)$, the entropy per oscillation $S(T)$, the specific heat capacity per oscillation $C_v(T)$ and the variance of position of particle vibration bound-state under anharmonics oscillator asymmetric potential system can be readily determined. As a result, all the thermodynamics anharmonics oscillator asymmetric potential can then be determined, provided that the macroscopic observable $S(T)$, $C_v(T)$ and the free energy are known as well. Subsequently the result of the classical thermodynamics are recovered from statistical mechanics. The contribution to the temperature dependent internal energy per oscillation is $U(T) = kT^2 \frac{\partial}{\partial T} (\ln(Z(T)))$. Taking the logarithm and the derivative yields in the partition function on the equation (11) for the temperature dependent internal energy per oscillation

$$U(T) = kT^2 \frac{\partial}{\partial T} (\ln(Z(T))) = kT^2 \left(\frac{1}{2} \frac{\partial}{\partial T} \ln \left(\frac{2\pi kT}{m\omega^2} \right) + \frac{\partial}{\partial T} \ln \left(1 + \frac{3bc}{m^2\omega^4} + \frac{b^2}{2mkT\omega^2} + \left(\frac{15c^2}{2m^3\omega^6} - \frac{d}{4m^2\omega^4} \right) (kT) + \frac{105d^2}{2m^4\omega^8} (kT)^2 \right) \right) \quad (12)$$

We can determine the temperature dependent specific heat capacity from the above by taking the derivative with respect to the temperature

$$C_v(T) = \frac{\partial}{\partial T} U(T) = \frac{k}{2} + 2 \left(\frac{15c^2}{2m^3\omega^6} - \frac{12d}{4m^2\omega^4} \right) (k^2T) + \frac{315d^2}{2m^4\omega^8} (k^3T^2) + \dots \quad (13)$$

Schematically sketch, the temperature dependent specific heat capacity in equation (13) has shape illustrate in figure (3). The temperature dependent entropy then defined as

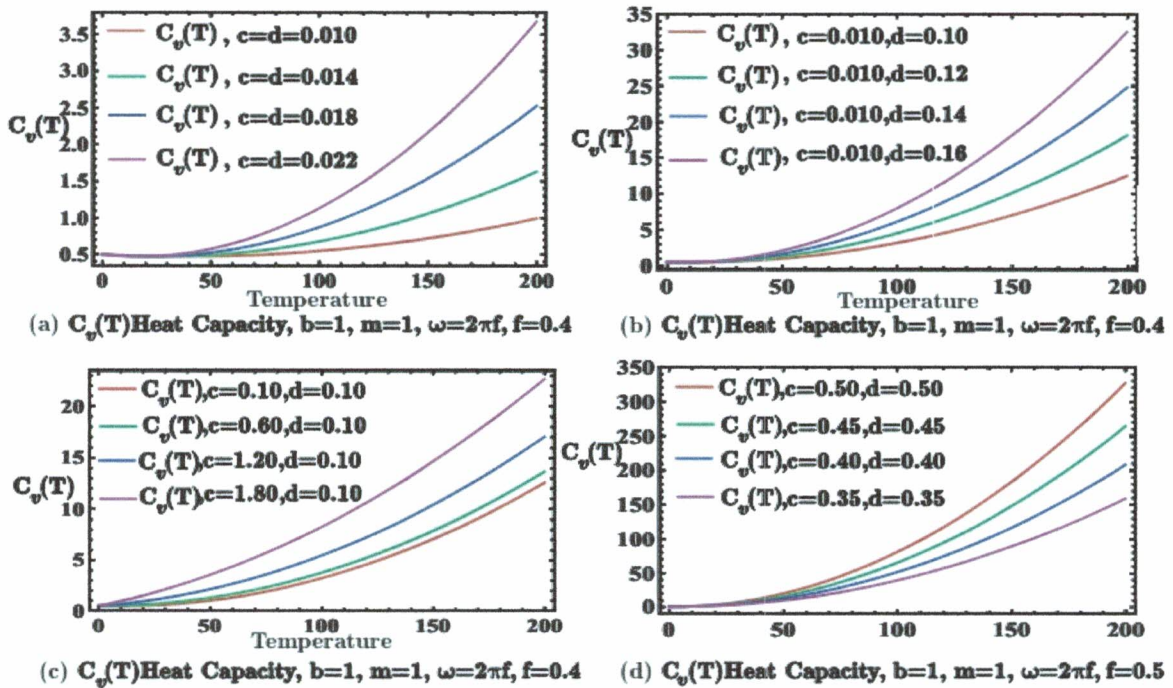


Fig. 3: Illustration showing a specific heat capacity of a particle moved bound by the anharmonics oscillator asymmetric potential and vary c and d parameter positive constant.

$$S(T) = k \ln(Z(T)) + \frac{U(T)}{T}$$

$$S(T) = \frac{k}{2} \ln\left(\frac{2\pi kT}{m\omega^2}\right) + k \ln\left(1 + \frac{3bc}{m^2\omega^4} + \frac{b^2}{2mkT\omega^2} + \left(\frac{15c^2}{2m^3\omega^6} - \frac{12d}{4m^2\omega^4}\right)(kT) + \frac{105d^2}{2m^4\omega^8}(kT)^2\right)$$

$$+ \frac{k}{2} - \frac{b^2}{2mkT\omega^2} + \left(\frac{15c^2}{2m^3\omega^6} - \frac{12d}{4m^2\omega^4}\right)(kT) + \frac{105d^2}{2m^4\omega^8}(kT)^2 + \dots \quad (14)$$

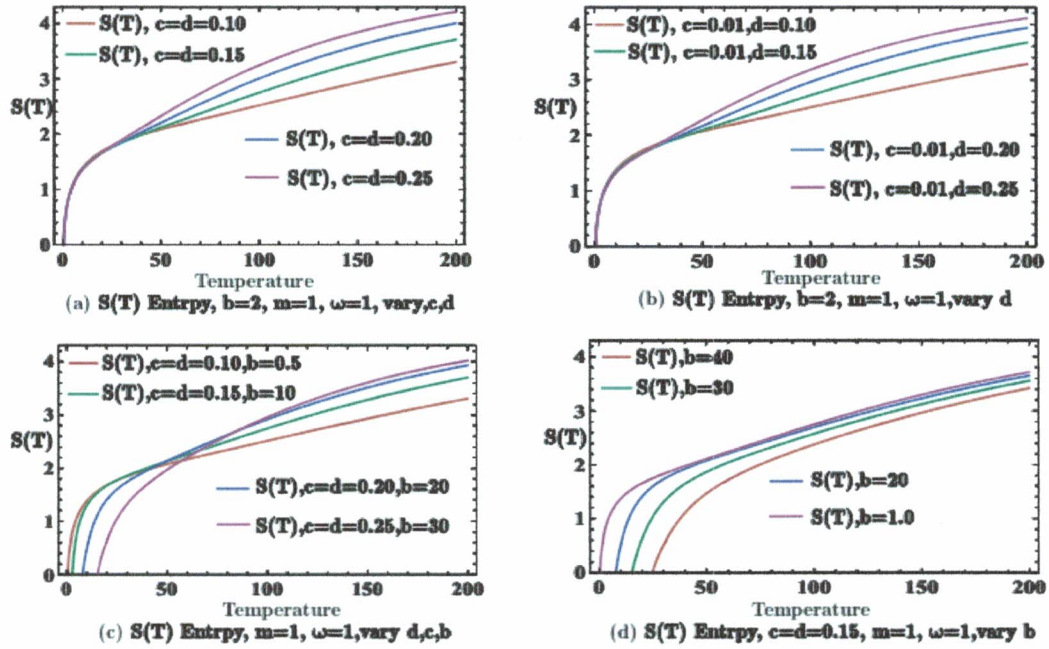


Fig. 4: Plot of the temperature-dependent entropy for particle move in the anharmonics oscillator asymmetric potential.

Sketch of the temperature dependent entropy for this system are given in figure (4). We can determine the temperature enthalpy from defined as $H(T) = U(T) + PV$ by setting P is the pressure, V is the volume of the anharmonics oscillator asymmetric potential system in $V(x)$

$$H(T) = \frac{kT}{2} - \frac{b^2}{2m\omega^2} + \left(\frac{b^2}{2m^3\omega^6} - \frac{12d^2}{4m^2\omega^4}\right)(kT)^2 + \frac{105d^2}{2m^4\omega^8}(kT)^3 + PV + \dots \quad (15)$$

The temperature dependent free energy of Helmholtz potential then defined in the anharmonics oscillator asymmetric potential system as $G(T) = U(T) - TS(T)$

$$G(T) = -\left(\frac{kT}{2} \ln\left(\frac{2\pi kT}{m\omega^2}\right) + kT \ln\left(1 + \frac{3bc}{m^2\omega^4} + \frac{b^2}{2m\omega^2 kT} + \left(\frac{15c^2}{2m^3\omega^6} - \frac{12d}{4m^2\omega^4}\right)(kT) + \frac{105d^2}{2m^4\omega^8}(kT)^2\right)\right) \quad (16)$$

The enthalpy and the temperature dependent free energy of Helmholtz potential of particle for the anharmonics oscillator asymmetric potential system is given by dependent on the parameter b , c , and d . Now consider the average value of the x coordinate for a single-particle oscillation one-dimensional system in the partition function of equation (11). The x coordinate takes on a continuous range of values, and the probability of observing the particle oscillator between x and $x + dx$ is $e^{-\beta V(x)} / Z(T)$ where β is $1/kT$. The summation over the infinitesimal probabilities is equivalent to an integration over the full range of x becomes

$$\langle x \rangle = \frac{1}{Z(T)} \int_{-\infty}^{\infty} x e^{-V(x)/kT} dx \quad (17)$$

Substituting the potential $V(x)$ and using the Taylor series expansion and using equation (3)-(7) into equation (17) result in the average value of position x dependent on temperature

$$\langle x \rangle = \frac{1}{Z(T)} \sqrt{\frac{2\pi kT}{m\omega^2}} \left\{ -\frac{b}{m\omega^2} - \frac{3c(kT)}{m^2\omega^4} + \frac{15bd(kT)}{m^3\omega^6} + \frac{105cd(kT)^2}{m^4\omega^8} + \dots \right\}. \quad (18)$$

The average value $\langle x^2 \rangle$ of the position of a particle is given by

$$\langle x^2 \rangle = \frac{1}{Z(T)} \int_{-\infty}^{\infty} x^2 e^{-V(x)/kT} dx \quad (19)$$

We again use the potential $V(x)$ and the Taylor series expansion and using equation (3)-(7) into equation (19), we get

$$\langle x^2 \rangle = \frac{1}{Z(T)} \sqrt{\frac{2\pi kT}{m\omega^2}} \left\{ \frac{3b^2}{2m^2\omega^4} + \left(\frac{1}{m\omega^2} + \frac{15c}{m^3\omega^6} \right) (kT) + \left(\frac{105c^2}{2m^4\omega^8} + \frac{15d}{m^3\omega^6} \right) (kT)^2 + \dots \right\} \quad (20)$$

The standard deviation $\sigma(T)$ of the position of the particle is defined as the root-mean-square deviation $\sigma(T) = \sqrt{\langle x^2 \rangle - \langle x \rangle^2}$ of the value of x from the average or the expectation value.

Therefore

$$\sigma(T) = \left(\frac{1}{Z(T)} \sqrt{\frac{2\pi kT}{m\omega^2}} \left\{ \frac{3b^2}{2m^2\omega^4} + \left(\frac{1}{m\omega^2} + \frac{15c}{m^3\omega^6} \right) (kT) + \left(\frac{105c^2}{2m^4\omega^8} + \frac{15d}{m^3\omega^6} \right) (kT)^2 + \dots \right\} - \left(\frac{1}{Z(T)} \sqrt{\frac{2\pi kT}{m\omega^2}} \left\{ -\frac{b}{m\omega^2} - \frac{3c(kT)}{m^2\omega^4} + \frac{15bd(kT)}{m^3\omega^6} + \frac{105cd(kT)^2}{m^4\omega^8} + \dots \right\} \right)^2 \right)^{\frac{1}{2}} \quad (21)$$

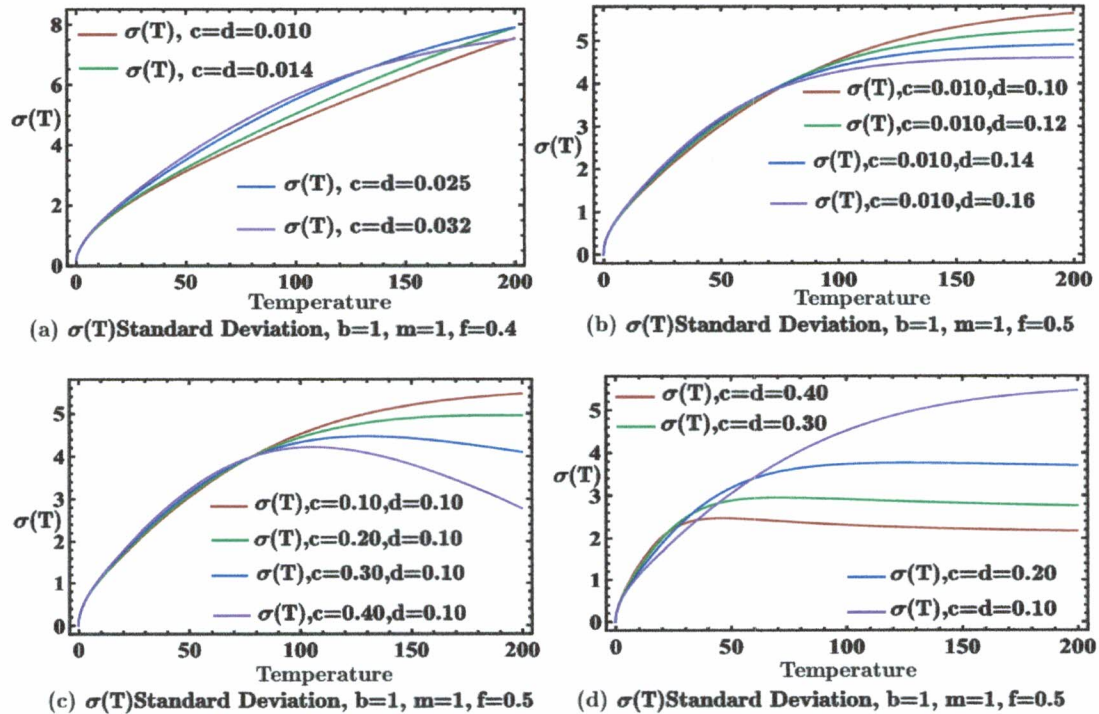


Fig. 5: Plot of the temperature-dependent standard deviation for particle move in the anharmonics oscillator asymmetric potential.

Discussion and Conclusion

Mention to figure (2), when we increase c and d parameter positive constant from 0.01 to 0.35 and 0.1 to 1.8. It will affect the partition function in the anharmonics oscillator asymmetric potential to be escalated. Mention to figure (3), when we increase c , d parameter positive constant from 0.01 to 0.5. It will affect the specific heat capacity in the anharmonics oscillator asymmetric potential to be increased. Mention to figure (4(a-d)), when we increase c, d parameter positive constant from 0.01 to 40. It will affect the entropy in the anharmonics oscillator asymmetric potential to be escalated. According to figure (5(a)), when we vary $c=d$ parameter positive constant from 0.01 to 0.032. It will affect the standard deviation in potential to be increased $\sigma(T)=8$. Summary, such the partition function, internal energy, specific heat capacity, entropy, enthalpy, free energy and standard deviation which depend on the c, d parameter and frequency (f).

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