รายชื่อวารสารทั้งหมด

พบวารสารทั้งหมด 1152 รายการ

*ท่านสามารถดูรายละเอียดของแต่ละวารสารได้โดยคลิกที่ชื่อของวารสาร

2774-0838

ISSN	E-ISSN	ชื่อไทย	ชื่ออังกฤษ	TCI กลุ่ มที่	สาขา	เว็บไซต์	нитвіна
2774- 0838	2774- 0757	<u>วารสารวิทยาศาสตร์และ</u> <u>เทคโนโลยี มหาวิทยาลัย</u> <u>ราชภัฎบุรีรัมย์</u>	Journal of Science and Technology Buriram Rajabhat University	2	Physical Sciences	https://ph02.tci- thaijo.org/index. php/scibru/inde x	ISSN และชื่อวารสาร (เดิม) คือ ISSN: : 1906-4713, วารสาร วิทยาศาสตร์และ เทคโนโลยี



ผลการประเมินคุณภาพวารสารที่อยู่ในฐานข้อมูล TCI

Show 10 v entries Search: 1906-4713

No.	Journal Name English	Journal Name Local	ISSN	E-ISSN	TCI Tier	Date for next submission
479	Journal of Science and Technology Buriram Rajabhat University	วารสารวิทยาศาสตร์และเทคโนโลยี มหาวิทยาลัยราชภัฏบุรีรัมย์	1906-4713	-	2	ไม่ก่อนวันที่ 1 ม.ค. 2568

Showing 1 to 1 of 1 entries (filtered from 951 total entries)

(htt ps:/ Thai-Journal Citation Index Centre (https://tci-thailand.org/) /so cial plu gins .lin e.m e/li neit /sh are url =ht tps %3 Α% 2F %2 Ftci thai lan d.or g% 2F %3 Fр %3 D37 96)

Manage consent

Previous

Next



วิทยาศาสตร์และเทคโนโลยี มหาวิทยาลัยราชภัฏบุรีรัมย์ ปีที่ 5 ฉบับที่ 2 เดือนกรกฎาคม - เดือนธันวาคม 2564

Journal of Science and Technology Buriram Rajabhat University

Vol. 5 No. 2

July – December 2021

ISSN 1906-4713

ที่ปรึกษา

รองศาสตราจารย์มาลิณี จุโทปะมา อธิการบดี มหาวิทยาลัยราชภัฏบุรีรัมย์
 ผู้ช่วยศาสตราจารย์ ดร.สุพัตรา รักการศิลป์ รองอธิการบดีฝ่ายวิชาการ มหาวิทยาลัยราชภัฏบุรีรัมย์
 ผู้ช่วยศาสตราจารย์ ดร.สุภาพร ชื่นชูจิตร คณบดีคณะวิทยาศาสตร์ มหาวิทยาลัยราชภัฏบุรีรัมย์

บรรณาธิการ

อาจารย์ ดร.ทิพวัลย์ แสนคำ

รองบรรณาธิการ

ผู้ช่วยศาสตราจารย์ ดร.สราวุธ แก้วศรี

กองบรรณาธิการ

1.	ศาสตราจารย์ ดร.ชิดชนก เหลือสินทรัพย์	จุฬาลงกรณ์มหาวิทยาลัย
2.	ศาสตราจารย์ ดร.อลงกลด แทนออมทอง	มหาวิทยาลัยขอนแก่น
3.	ศาสตราจารย์ ดร.ปรีชา ประเทพา	มหาวิทยาลัยมหาสารคาม
4.	ศาสตราจารย์ ดร.ไพโรจน์ ประมวล	มหาวิทยาลัยมหาสารคาม
5.	รองศาสตราจารย์ ดร.วิภาวี หีบแก้ว	มหาวิทยาลัยขอนแก่น
6.	ผู้ช่วยศาสตราจารย์ ดร.เกตุจันทร์ จำปาไชยศรี	มหาวิทยาลัยนเรศวร
7.	ผู้ช่วยศาสตราจารย์ ดร.เสรี ปานซาง	มหาวิทยาลัยราชภัฏเชียงใหม่
8.	ผู้ช่วยศาสตราจารย์ น.สพ.ประยุทธ กุศลรัตน์	มหาวิทยาลัยราชภัฏนครราชสีมา
9.	ผู้ช่วยศาสตราจารย์ ดร.ไชยยันต์ สกุลไทย	มหาวิทยาลัยราชภัฏมหาสารคาม
10.	อาจารย์ ดร.สุนันทา กลิ่นถาวร	มหาวิทยาลัยราชภัฏมหาสารคาม
11.	ผู้ช่วยศาสตราจารย์ ดร.ถิรนันท์ สอนแก้ว	มหาวิทยาลัยราชภัฏลำปาง
12.	อาจารย์ ดร.ชาติวุฒิ ธนาจิรันธร	มหาวิทยาลัยราชภัฏบุรีรัมย์

ผู้ทรงคุณวุฒิประจำฉบับ

1.	รองศาสตราจารย์	์ ดร.ชัยณรงค์	้อธิสกล
Τ.	00 111 101710 10 100	, 110.00000 all	00011101

2. รองศาสตราจารย์ ดร.จรัส สว่างทัพ

3. ผู้ช่วยศาสตราจารย์ ดร.ชนินทร์ มหัทธนชัย

4. ผู้ช่วยศาสตราจารย์ ดร.ณัฐวุฒิ โสมะเกษตรินทร

5. ผู้ช่วยศาสตราจารย์ ดร.ธรัช อารีราษฎร์

6. ผู้ช่วยศาสตราจารย์ ดร.นพปฎล เสงี่ยมศักดิ์

7. ผู้ช่วยศาสตราจารย์ ดร.เมตตา เถาว์ชาลี

8. ผู้ช่วยศาสตราจารย์ ดร.ราตรี คำโมง

9. ผู้ช่วยศาสตราจารย์ ดร.วรปภา อารีราษฎร์

10. ผู้ช่วยศาสตราจารย์ ดร.วิฑูรย์ พรมมี

11. ผู้ช่วยศาสตราจารย์ ดร.เสริมศักดิ์ พงษ์เมษา

12. ผู้ช่วยศาสตราจารย์ ดร.เสรี ปานซาง

13. ผู้ช่วยศาสตราจารย์ ดร.โสภา แคนสี

14. ผู้ช่วยศาสตราจารย์ น.สพ.ประยุทธุ กุศลรัตน์

15. ผู้ช่วยศาสตราจารย์ น.สพ.สมพงศ์ วงษ์มา

16. ผู้ช่วยศาสตราจารย์ ดร.เดือนฉาย ไชยบุตร

17. ผู้ช่วยศาสตราจารย์ ดร.กมลรัตน์ สมใจ

18. ผู้ช่วยศาสตราจารย์ ดร.จาริณี ม้าแก้ว

19. ผู้ช่วยศาสตราจารย์ ดร.วรวัฒน์ พรหมเด่น

20. ผู้ช่วยศาสตราจารย์ ดร.วิไลรัตน์ ยาทองไชย

21. ผู้ช่วยศาสตราจารย์ ดร.สมศักดิ์ จีวัฒนา

22. ดร.วิริญรัชญ์ สื่อออก

มหาวิทยาลัยเกษตรศาสตร์ มหาวิทยาลัยราชภัฏบุรีรัมย์ มหาวิทยาลัยราชภัฏเชียงใหม่ มหาวิทยาลัยเทคโนโลยีราชมงคลธัญบุรี มหาวิทยาลัยราชภัฏมหาสารคาม มหาวิทยาลัยมหาสารคาม มหาวิทยาลัยราชภัฏมหาสารคาม มหาวิทยาลัยราชภัฏอุตรดิตถ์ มหาวิทยาลัยราชภัฏมหาสารคาม มหาวิทยาลัยเทคโนโลยีราชมงคลล้านนา มหาวิทยาลัยราชภัฏเชียงใหม่ มหาวิทยาลัยราชภัฏเชียงใหม่ มหาวิทยาลัยมหาสารคาม มหาวิทยาลัยราชภัฏนครราชสีมา มหาวิทยาลัยราชภัฏนครราชสีมา มหาวิทยาลัยราชภัฏเพชรบูรณ์ มหาวิทยาลัยราชภัฏบุรีรัมย์ มหาวิทยาลัยราชภัฏบุรีรัมย์ มหาวิทยาลัยราชภัฏบุรีรัมย์ มหาวิทยาลัยราชภัฏบุรีรัมย์

มหาวิทยาลัยราชภัฏบุรีรัมย์

มหาวิทยาลัยราชภัฏบุรีรัมย์

สารบัญ

เรื่อง
การพัฒนาระบบบริหารจัดการร้านอิเล็กทรอนิกส์เซอร์วิส
Evaluation of the Time-dependent Electric Charge in the Series of RLC Circuit Loop Under the Time-dependent Voltage Cosine Function
เครื่องสลัดน้ำมันพริกทอดเพื่อเพิ่มคุณภาพผลิตภัณฑ์ : กรณีศึกษาชุมชนบ้านหนองบัวแดง
การพัฒนาแผ่นฝ้าเพดานจากใยไม้ผสมผงยิปซัมที่มีคุณสมบัติเป็นฉนวนป้องกันความร้อน45 อธิพันธ์ ลอยเมืองกลาง และ อัคร เสมรบุณย์
การพัฒนาผลิตภัณฑ์ขัดผิวจากเศษกะลาแมคคาเดเมีย สำหรับวิสาหกิจชุมชนกลุ่มผู้ผลิต
การประยุกต์ใช้เทคโนโลยีดิจิทัลเพื่อถ่ายทอดองค์ความรู้ภูมิปัญญาการผลิตผ้าทอใยกัญชง
ผลการเสริมโพรไบโอติกแบซิลลัส ซับทิลิส (Bacillus subtilis) ในอาหารต่อสมรรถภาพ การผลิต ของไก่พื้นเมืองพันธุ์เหลืองหางขาว
การพัฒนาระบบพาณิชย์อิเล็กทรอนิกส์สำหรับสวนไผ่ทองสุข ต.โคกตูม อ.เมือง จังหวัดลพบุรี101 ดวงพร ไม้ประเสริฐ

Evaluation of the Time-dependent Electric Charge in the Series of RLC Circuit Loop Under the Time-dependent Voltage Cosine Function

Nutnicha Masoongnoen², Sanit Suwanwong¹, Treenuch Ellis¹,

and Artit Hutem1*

Physics Division, Faculty of Science and Technology, Phetchabun Rajabhat University Thailand¹

Wittayanukulnaree School, Phetchabun, Thailand ²

Email: artithutem@pcru.ac.th

Received: July 27, 2021

Revised: September 1, 2021

Accepted: September 14, 2021

Abstract

In this work, we developed a model of mathematics and physics for the series of

the RLC circuit loop. The purpose of this study is to evaluate for finding the time-

dependent electric charge that is a consequence of time-dependent voltage force.

Which the voltage force is in the cosine function. We calculated by using the second-order

non-homogeneous ordinary differential equation and integration by part technique. We

can find that the time-dependent electric charge corresponds to capacitance but is

inversely proportional to induction. The time-dependent electric charge is in contrast to

the charge. If δ and σ have slightly different values the time-dependent electric charge

behaves like an oscillation wave group.

Keyword: Time-dependent Electric Charge, Time-dependent Voltage Force, Wave Group

INTRODUCTION

In one loop of RLC circuit or second-order circuit is from the including of resistor,

inductor and capacitor where in this case, we analyze where the voltage depends on

time and it is in cosine function. Mohazzab J. et al. (2008) study RLC circuit response

and analysis using state space method. So he easily fined the response and stability of

the RLC circuit or second-order circuit and also with the help of the response of the

วารสารวิทยาศาสตร์และเทคโนโลยี ปีที่ 5 ฉบับที่ 2

RLC is examined from different input functions by using Matlab. The analysis of an RLC circuit becomes too simpler. Sonam (2015) can use Matlab for analyzing the free natural angular frequency repercussion of the second-order circuit, time repercussion of the circuit. To analyse other standard second-order circuit conformation such as low pass and high pass RLC network, we use interactive GUI. In GUI you can change the RLC or second-order circuit parameter and see the fructification on the time and free natural angular frequency repercussion in real time.

Dino (2019) study a low frequency hook with positive imaginary part in the impedance spectrum can be explained by several phenomena and it occurs in a number of photo-electrochemical system. Ahammodullah (2019) have successfully applied the Kirchhoff's voltage rule (KVL) modified into non-homogeneous second-order differential equation to series circuits loop containing an electromotive force, resistor, inductor and capacitor. The purpose of this paper is solving the series of RLC circuits or second-order circuit (Sokol $et\ al.$, 2013) in the driven force dependent on time in cosine function form by using Wronskian's method form of the non-homogeneous second-order differential equation. As the result gives 2 cases of charge parameter that the electric charge value corresponds to the capacitance but inversely proportional to induction, anywhere if the value of the parameter δ is defined as ratio between resistance with double inductance and the free natural angular frequency (σ) have more different values, electric charge would have the behavior of themselves like wave group.

MATERIALS AND METHODS

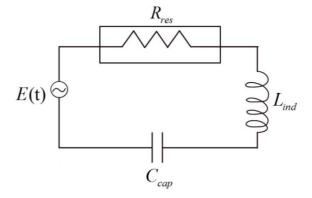


Figure 1: Representation the series of RLC circuit with time-dependent voltage.

The electrical system was made up of a capacitor $(C_{\it Cap})$, resistance $(R_{\it res})$, inductance $(L_{\it ind})$ and applied voltage force (E(t)) is one of the most important resonate systems in figure 1 (Kishore, 2008). In this work-study, the voltage force depends on time and in cosine form. Thus the sum of the voltages encountered in going around the whole circuit must be zero. Let us analyze the circuit of this case by applying Kirchhoff's rule, we get

$$E(t) = E_R + E_L + E_C \tag{1}$$

From the article above, the time-dependent applied voltage force is cosine function form $E(t) = E_0 e^{-\eta^2 t} \cos^2(\sigma t)$ (Hutem & Masoongnoen,. 2021) which E_0 is the initial voltage force, η is the coefficient of damping of the voltages force and σ is a constant and is called the free natural angular frequency. After substituting E(t) into Equation (1), we obtain

$$E_0 e^{-\eta^2 t} \cos^2(\sigma t) = E_0 + E_1 + E_2 \tag{2}$$

We may write the voltage drop across the resistance $E_R = I_{sys}R_{res}$, across the capacitor it is $E_C = q/C_{cap}$ and across the inductance it is $E_L = L_{ind}\left(dI_{sys}/dt\right)$ (Atamp , 1990) and in the new form where $I_{sys} = dq/dt$ (Goldstein and Safko, 2002), we get

$$\frac{d^2q}{dt^2} + \frac{R_{res}}{L_{ind}} \frac{dq}{dt} + \frac{1}{L_{ind}C_{con}} q = \frac{E_0}{L_{ind}} e^{-\eta^2 t} \cos^2(\sigma t)$$
 (3)

Therefore, we can set the new form where R_{res}/L_{ind} is defined as 2δ , $1/L_{ind}C_{cap}$ is defined as σ^2 and E_0/L_{ind} is defined as ε_0 . Figure 1. represents an analogous electrical driven oscillator with an applied voltage force emf source given by $E(t)=E_0e^{-\eta^2t}\cos^2(\sigma t)$. We can rewrite Equation (3) as

$$\frac{d^2q}{dt^2} + 2\delta \frac{dq}{dt} + \sigma^2 q = \varepsilon_0 e^{-\eta^2 t} \cos^2(\sigma t)$$
 (4)

Equation 4 is called a non-homogeneous second-order differential equation (Tikjha *et al.*, 2018) that can give the solution by using the summation of two part as

$$q(t) = q_{\scriptscriptstyle C}(t) + q_{\scriptscriptstyle P}(t) \,, \tag{5}$$

where $q_{\scriptscriptstyle P}(t)$ is the time-dependent electric charge particular solution of an inhomogeneous differential equation and the time-dependent electric charge of complementary function $q_c(t)$ is solution of the corresponding homogeneous differential equation (that is equation (4) with the right side equal to zero number), then equation (5) is also a solution of the non-homogeneous second-order differential equation. $q_c(t)$ is the time-dependent electric charge of complementary function solution of the homogeneous second-order differential equation (Riley & Hobson , 2006)

$$\frac{d^2q}{dt^2} + 2\delta \frac{dq}{dt} + \sigma^2 q = 0 \tag{6}$$

Let us solve the auxiliary equation as $m^2+2\delta\,m+\sigma^2=0$. The auxiliary equation has the roots

$$m = -\delta \pm \sqrt{\delta^2 - \sigma^2} \tag{7}$$

For this case, we can select to analyze the underdamped for the convenience to substitution where $\delta < \sigma$. We must have $m = -\delta \pm i \hbar$ and $\hbar = \sqrt{\sigma^2 - \delta^2}$. Thus, the complementary function solution of the homogeneous differential equation of Equation (6) is

$$q_C(t) = e^{-\delta t} \left(\alpha \sin(\lambda t) + \beta \cos(\lambda t) \right).$$
 (8)

We can find the value of α is a constant and β is a constant by setting the boundary conditions $q(0)=q_0$ (the initial electric charge) and dq(0)/dt=0. Thus, we consider the charge depends on time when t=0 as

$$q_{C}(0) = 0 + \beta \cos(0), \Rightarrow \beta = q_{0}$$
(9)

where $\beta=q_0$ is the initial charge. Next, we can find the derivative of charge depends on time term as

$$\frac{dq_{c}(t)}{dt} = e^{-\delta t} \left(\alpha \lambda \cos(\lambda t) - \beta \sin(\lambda t) \right) - \delta e^{-\delta t} \left(\alpha \lambda \sin(\lambda t) + \beta \cos(\lambda t) \right)$$

From the initial condition of derivative part, we must have

$$\alpha = \frac{\delta q_0}{\hbar} = \frac{\delta q_0}{\sqrt{\sigma^2 - \delta^2}},\tag{10}$$

where $R_{res}/2L_{ind}$ is defined as δ , q_0 is the initial electric charge. Thus, the homogeneous of Equation (6) gives the complementary function solution as

$$q_C(t) = e^{-\delta t} \left(\frac{\delta q_0}{\hbar} \sin(\hbar t) + q_0 \cos(\hbar t) \right)$$
 (11)

Then, let us seek the particular solution of Equation (5) by using Wronskian method (Susan, 2004). We try the trigonometry and exponential solution of the form that

$$\begin{aligned} y_1(t) &= e^{-\delta t} \sin(\lambda t) \quad , & y_2(t) &= e^{-\delta t} \cos(\lambda t) \\ y_1'(t) &= \lambda e^{-\delta t} \cos(\lambda t) - \delta e^{-\delta t} \sin(\lambda t) & y_2'(t) &= -\lambda e^{-\delta t} \sin(\lambda t) - \delta e^{-\delta t} \cos(\lambda t) \end{aligned}$$

We can find the solution of Wronskian method by the form below

$$W_{Wr} = \begin{vmatrix} y_1(t) & y_2(t) \\ y_1'(t) & y_2'(t) \end{vmatrix}$$

Thus, we get

We get
$$W_{Wr} = \begin{vmatrix} e^{-\delta t} \sin(\hat{\lambda}t) & e^{-\delta t} \cos(\hat{\lambda}t) \\ \hat{\lambda}e^{-\delta t} \cos(\hat{\lambda}t) - \delta e^{-\delta t} \sin(\hat{\lambda}t) & -\hat{\lambda}e^{-\delta t} \sin(\hat{\lambda}t) - \delta e^{-\delta t} \cos(\hat{\lambda}t) \end{vmatrix}$$

$$W_{Wr} = -\hat{\lambda}e^{-2\delta t}$$
(12)

Consider the first Wronskian method and we satisfy the solution as

$$W_{W1} = \begin{vmatrix} 0 & e^{-\delta t} \cos(\hbar t) \\ \varepsilon_0 e^{-\eta^2 t} \cos^2(\sigma t) & -(\hbar e^{-\delta t} \sin(\hbar t) + \delta e^{-\delta t} \cos(\hbar t)) \end{vmatrix}$$

$$W_{W1} = -\varepsilon_0 e^{-(\eta^2 + \delta)t} \cos^2(\sigma t) \cos(\hbar t)$$
(13)

Consider the second Wronskian method and we satisfy the solution as

$$W_{W2} = \begin{vmatrix} e^{-\delta t} \sin(\hbar t) & 0\\ \hbar e^{-\delta t} \cos(\hbar t) - \delta e^{-\delta t} \sin(\hbar t) & \varepsilon_0 e^{-\eta^2 t} \cos^2(\sigma t) \end{vmatrix}$$

$$W_{W2} = \varepsilon_0 e^{-(\eta^2 + \delta)t} \cos^2(\sigma t) \sin(\hbar t) \tag{14}$$

We can satisfy the derivative equation for Wronskian method (Sadri ,1991) as

$$\frac{du_1}{dt} = \frac{W_{W1}}{W_{Wc}} = \frac{-\varepsilon_0 e^{-(\eta^2 + \delta)t} \cos^2(\sigma t) \cos(\hat{\lambda}t)}{-\hat{\lambda}e^{-2\delta t}} = \frac{\varepsilon_0}{\hat{\lambda}} e^{(\delta - \eta^2)t} \cos(\hat{\lambda}t) \cos^2(\sigma t)$$

So we set the parameter of $\chi = (\delta - \eta^2)$ as

$$\frac{du_1}{dt} = \frac{\varepsilon_0}{\lambda} e^{\chi t} \cos^2(\sigma t) \cos(\lambda t)$$

$$\frac{du_1}{dt} = \frac{\varepsilon_0}{\lambda} e^{\chi t} \left(\frac{1 + \cos(2\sigma t)}{2}\right) \cos(\lambda t)$$

$$u_1 = \frac{\varepsilon_0}{2\lambda} \int e^{\chi t} \left(1 + \cos(2\sigma t)\right) \cos(\lambda t) dt \tag{15}$$

Thus,

$$\therefore u_1(t) = \frac{\varepsilon_0}{2\hbar} \left(\int e^{\chi t} \cos(\hbar t) dt + \int e^{\chi t} \cos(\hbar t) \cos(\sigma 2t) dt \right)$$
 (16)

We can find the first integral function by using integration by part technique for estimation the right hand of Equation (16).

$$\int e^{\chi t} \cos(\lambda t) dt = e^{\chi t} \left(\frac{\lambda \sin(\lambda t) + \chi \cos(\lambda t)}{(\lambda^2 + \chi^2)} \right)$$
 (17)

Then consider the second integral function by using integration by part technique for estimation the right-hand side of Equation (16).

$$\int e^{\chi t} \cos(\lambda t) \cos(2\sigma t) dt = \frac{1}{2} \left[\frac{e^{\chi t} \left((2\sigma + \lambda) \sin((2\sigma + \lambda)t) + \chi \cos((2\sigma + \lambda)t) \right)}{((2\sigma + \lambda)^2 + \chi^2)} + \frac{e^{\chi t} \left((\lambda - 2\sigma) \sin((\lambda - 2\sigma)t) + \chi \cos((\lambda - 2\sigma)t) \right)}{((\lambda - 2\sigma)^2 + \chi^2)} \right]$$
(18)

We will get the solution of derivative equation for Wronskian method as

$$u_{1}(t) = \frac{\varepsilon_{0}e^{\chi t}}{2\hbar} \left(\frac{\left(\hbar \sin(\hbar t) + \chi \cos(\hbar t)\right)}{(\hbar^{2} + \chi^{2})} + \frac{\left((2\sigma + \hbar)\sin((2\sigma + \hbar)t) + \chi \cos((2\sigma + \hbar)t)\right)}{2((2\sigma + \hbar)^{2} + \chi^{2})} + \frac{(\hbar - 2\sigma)\sin((\hbar - 2\sigma)t) + \chi \cos((\hbar - 2\sigma)t)}{2((\hbar - 2\sigma)^{2} + \chi^{2})} \right)$$

$$(19)$$

The derivative for the second Wronskian method can be satisfied as

$$\frac{du_2}{dt} = \frac{W_2}{W} = -\frac{\varepsilon_0 e^{-(\eta^2 + \delta)t} \cos^2(\sigma t) \sin(\lambda t)}{\lambda e^{-2\delta t}} = -\frac{\varepsilon_0}{2\lambda} e^{\lambda t} (1 + \cos(2\sigma t)) \sin(\lambda t)$$

Therefore.

$$u_2(t) = -\frac{\varepsilon_0}{2\hbar} \left(\int e^{\lambda t} \sin(\hbar t) dt + \int e^{\lambda t} \sin(\hbar t) \cos(2\sigma t) dt \right)$$
 (20)

We can find the first derivative in the right hand of Equation (20) by using integration by part technique for estimation.

$$\int e^{\chi t} \sin(\lambda t) dt = \frac{e^{\chi t} \left(\chi \sin(\lambda t) - \lambda \cos(\lambda t) \right)}{(\chi^2 + \lambda^2)}$$
 (21)

Then, we also use integration by part technique for estimate the second integral function in the right hand of Equation (20) as

$$\int e^{\chi t} \cos(2\sigma t) \sin(\lambda t) dt = \frac{e^{\chi t} \left(\chi \sin((\lambda - 2\sigma)t) - (\lambda - 2\sigma)\cos((\lambda - 2\sigma)t) \right)}{2(\chi^2 + (\lambda - 2\sigma)^2)} + \frac{e^{\chi t} \left(\chi \sin((\lambda + 2\sigma)t) - (\lambda + 2\sigma)\cos((\lambda + 2\sigma)t) \right)}{2(\chi^2 + (\lambda + 2\sigma)^2)}$$
(22)

Thus, we get the solution of derivative of the second Wronskian method as

$$u_{2}(t) = -\frac{\varepsilon_{0}e^{xt}}{2\lambda} \left(\frac{\left(\chi \sin(\lambda t) - \lambda \cos(\lambda t)\right)}{(\chi^{2} + \lambda^{2})} + \frac{\left(\chi \sin((\lambda - 2\sigma)t) - (\lambda - 2\sigma)\cos((\lambda - 2\sigma)t)\right)}{2(\chi^{2} + (\lambda - 2\sigma)^{2})} + \frac{\left(\chi \sin((\lambda + 2\sigma)t) - (\lambda + 2\sigma)\cos((\lambda + 2\sigma)t)\right)}{2(\chi^{2} + (\lambda + 2\sigma)^{2})} \right)$$

$$(23)$$

The particular of a non-homogeneous, the second order, linear equation in Equation (4) with the constant ${\it C}$ is

$$\begin{split} q_{p}(t) &= \frac{\varepsilon_{0}e^{-\eta^{2}t}}{2\hbar}\sin(\hbar t) \left(\frac{\left(\hbar \sin(\hbar t) + (\delta - \eta^{2})\cos(\hbar t)\right)}{(\hbar^{2} + (\delta - \eta^{2})^{2})} + \frac{\left((2\sigma + \hbar)\sin((2\sigma + \hbar)t) + (\delta - \eta^{2})\cos((2\sigma + \hbar)t)\right)}{2((2\sigma + \hbar)^{2} + (\delta - \eta^{2})^{2})} \right. \\ &\quad + \frac{(\hbar - 2\sigma)\sin((\hbar - 2\sigma)t) + (\delta - \eta^{2})\cos((\hbar - 2\sigma)t)}{2((\hbar - 2\sigma)^{2} + (\delta - \eta^{2})^{2})} - \frac{\varepsilon_{0}e^{-\eta^{2}t}}{2\hbar}\cos(\hbar t) \left(\frac{\left((\delta - \eta^{2})\sin(\hbar t) - \hbar \cos(\hbar t)\right)}{((\delta - \eta^{2})^{2} + \hbar^{2})} \right. \\ &\quad + \frac{\left((\delta - \eta^{2})\sin((\hbar - 2\sigma)t) - (\hbar - 2\sigma)\cos((\hbar - 2\sigma)t)\right)}{2((\delta - \eta^{2})^{2} + (\hbar - 2\sigma)^{2})} + \frac{\left((\delta - \eta^{2})\sin((\hbar + 2\sigma)t) - (\hbar + 2\sigma)\cos((\hbar + 2\sigma)t)\right)}{2((\delta - \eta^{2})^{2} + (\hbar + 2\sigma)^{2})} + C \end{split}$$

We can find the parameter of C by using the initial condition as $q_{_{p}}(0) = q_{_{0}}$.

$$\begin{split} q_p(0) &= -\frac{\varepsilon_0}{2\hbar} \Biggl(-\frac{\hbar}{(\hbar^2 + (\delta - \eta^2)^2)} - \frac{(\hbar - 2\sigma)}{2((\hbar - 2\sigma)^2 + (\delta - \eta^2)^2)} - \frac{(\hbar + 2\sigma)}{2((\hbar + 2\sigma)^2 + (\delta - \eta^2)^2)} \Biggr) + C \\ C &= q_0 - \frac{\varepsilon_0}{2\hbar} \Biggl(\frac{\hbar}{(\hbar^2 + (\delta - \eta^2)^2)} + \frac{(\hbar - 2\sigma)}{2((\hbar - 2\sigma)^2 + (\delta - \eta^2)^2)} + \frac{(\hbar + 2\sigma)}{2((\hbar + 2\sigma)^2 + (\delta - \eta^2)^2)} \Biggr) \end{split} \tag{25}$$

Therefore, we get the particular solution of a non-homogeneous, second order, differential equation as

$$\begin{split} q_{p}(t) &= \frac{\varepsilon_{0}e^{-\eta^{2}t}}{2\hbar}\sin(\hbar t) \left(\frac{\left(\hbar \sin(\hbar t) + (\delta - \eta^{2})\cos(\hbar t)\right)}{(\hbar^{2} + (\delta - \eta^{2})^{2})} + \frac{\left((2\sigma + \hbar)\sin((2\sigma + \hbar)t) + (\delta - \eta^{2})\cos((2\sigma + \hbar)t)\right)}{2((2\sigma + \hbar)^{2} + (\delta - \eta^{2})^{2})} \right. \\ &+ \frac{(\hbar - 2\sigma)\sin((\hbar - 2\sigma)t) + (\delta - \eta^{2})\cos((\hbar - 2\sigma)t)}{2((\hbar - 2\sigma)^{2} + (\delta - \eta^{2})^{2})} - \frac{\varepsilon_{0}e^{-\eta^{2}t}}{2\hbar}\cos(\hbar t) \left(\frac{\left((\delta - \eta^{2})\sin(\hbar t) - \hbar\cos(\hbar t)\right)}{((\delta - \eta^{2})^{2} + \hbar^{2})} \right. \\ &+ \frac{\left((\delta - \eta^{2})\sin((\hbar - 2\sigma)t) - (\hbar - 2\sigma)\cos((\hbar - 2\sigma)t)\right)}{2((\delta - \eta^{2})^{2} + (\hbar - 2\sigma)^{2})} + \frac{\left((\delta - \eta^{2})\sin((\hbar + 2\sigma)t) - (\hbar + 2\sigma)\cos((\hbar + 2\sigma)t)\right)}{2((\delta - \eta^{2})^{2} + (\hbar + 2\sigma)^{2})} \\ &+ q_{0} - \frac{\varepsilon_{0}}{2\hbar} \left(\frac{\hbar}{(\hbar^{2} + (\delta - \eta^{2})^{2})} + \frac{(\hbar - 2\sigma)}{2((\hbar - 2\sigma)^{2} + (\delta - \eta^{2})^{2})} + \frac{(\hbar + 2\sigma)}{2((\hbar - 2\sigma)^{2} + (\delta - \eta^{2})^{2})} \right) \end{split}$$

Thus, we get the parameter of the time-dependent electric charge in Equation (5) as

$$\begin{split} q(t) &= e^{-\delta t} \left(\frac{\delta q_0}{\hbar} \sin(\hbar t) + q_0 \cos(\hbar t) \right) + \frac{\varepsilon_0 e^{-\eta^2 t}}{2\hbar} \sin(\hbar t) \left(\frac{\left(\hbar \sin(\hbar t) + (\delta - \eta^2) \cos(\hbar t) \right)}{(\hbar^2 + (\delta - \eta^2)^2)} \right. \\ &\quad \left. + \frac{\left((2\sigma + \hbar) \sin((2\sigma + \hbar)t) + (\delta - \eta^2) \cos((2\sigma + \hbar)t) \right)}{2((2\sigma + \hbar)^2 + (\delta - \eta^2)^2)} + \frac{(\hbar - 2\sigma) \sin((\hbar - 2\sigma)t) + (\delta - \eta^2) \cos((\hbar - 2\sigma)t)}{2((\hbar - 2\sigma)^2 + (\delta - \eta^2)^2)} \right) \end{split}$$

$$-\frac{\varepsilon_{0}e^{-\eta^{2}t}}{2\hbar}\cos(\hbar t)\left(\frac{\left((\delta-\eta^{2})\sin(\hbar t)-\hbar\cos(\hbar t)\right)}{\left((\delta-\eta^{2})^{2}+\hbar^{2}\right)}+\frac{\left((\delta-\eta^{2})\sin((\hbar-2\sigma)t)-(\hbar-2\sigma)\cos((\hbar-2\sigma)t)\right)}{2\left((\delta-\eta^{2})^{2}+(\hbar-2\sigma)^{2}\right)}+\frac{\left((\delta-\eta^{2})\sin((\hbar-2\sigma)t)-(\hbar-2\sigma)\cos((\hbar-2\sigma)t)\right)}{2\left((\delta-\eta^{2})^{2}+(\hbar-2\sigma)^{2}\right)}+\frac{\left((\delta-\eta^{2})\sin((\hbar+2\sigma)t)-(\hbar+2\sigma)\cos((\hbar+2\sigma)t)\right)}{2\left((\delta-\eta^{2})^{2}+(\hbar+2\sigma)^{2}\right)}+q_{0}-\frac{\varepsilon_{0}}{2\hbar}\left(\frac{\hbar}{(\hbar^{2}+(\delta-\eta^{2})^{2})}+\frac{(\hbar-2\sigma)}{2\left((\hbar-2\sigma)^{2}+(\delta-\eta^{2})^{2}\right)}+\frac{(\hbar+2\sigma)}{2\left((\hbar+2\sigma)^{2}+(\delta-\eta^{2})^{2}\right)}\right)$$

$$(27)$$

From Equation (27) is the charge depends on time show in program. Putting this into program Mathematica for plotting graph.

Case1:the time-dependent electric charge where δ and σ have more different value. ($\sigma\gg\delta$)

The time-dependent electric charge can behave like the underdamp wave. (Teoh & Rahifa, 2018).

Case2: the time-dependent electric charge where δ and σ have slightly different value. $(\sigma > \delta)$

The time-dependent electric charge can behave like the wave group.

RESULTS

We can explain of numerical and result of the time-dependent of electric charge in Equation (27) which effects by cosine voltage force as time-dependent charge as figure (2), figure (3), figure (4) and figure (5)

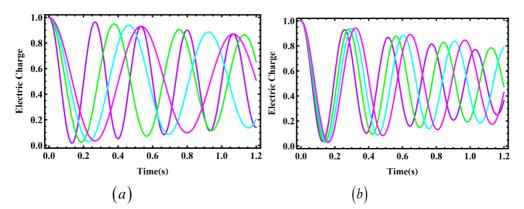


Figure 2: Illustration showing the relation parameter between charges depends on time where δ and σ have more different value. (a) Representation the relation between the charge and time when the capacitor value is unable (The purple solid

line is $C_{Cap1}=4\mu F$. The green solid line is $C_{Cap2}=8\mu F$. The light blue solid line is $C_{Cap3}=12\mu F$ The pink solid line is $C_{Cap4}=16\mu F$), (b) Representation the relation between the charge and time when the Inductor value is unable (The purple solid line is $L_{ind1}=210H$. The green solid line is $L_{ind2}=250H$. The light blue solid line is $L_{ind3}=290H$ The pink solid line is $L_{ind4}=330H$).

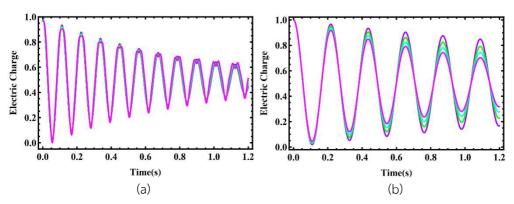


Figure 3: Illustration showing the relation parameter between charges depends on time where δ and σ have more different value. (a) Representation the relation between the electric charge and time when the initial applied voltage force (E_0) value is unable (The purple solid line is $E_{01}=0.1\,Vol$). The green solid line is $E_{02}=8\,kV$). The light blue solid line is $E_{03}=16\,kV$. The pink solid line is $E_{04}=26\,kV$), (b) Representation the relation between the electric charge and time when the resistance (R_{res}) value is unable (The purple solid line is $R_{res1}=100\,W$). The green solid line is $R_{res3}=150\,W$).

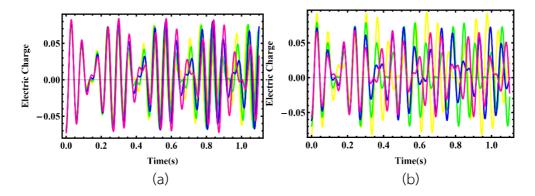


Figure 4: Illustration showing the relation parameter between charges depends on time where δ and σ have slightly different value $(\sigma > \delta)$. (a) Representation the relation between the charge and time when the capacitor value is unable (The yellow solid line is $C_{cap1} = 410 \mu F$. The green solid line is $C_{cap2} = 412 \mu F$. The blue solid line is $C_{cap3} = 414 \mu F$ The pink solid line is $C_{cap4} = 416 \mu F$), (b) Representation the relation between the charge and time when the inductor value is unable (The yellow solid line is $L_{ind1} = 1.10 H$). The green solid line is $L_{ind2} = 1.12 H$. The blue solid line is $L_{ind3} = 1.14 H$ The pink solid line is $L_{ind4} = 1.16 H$).

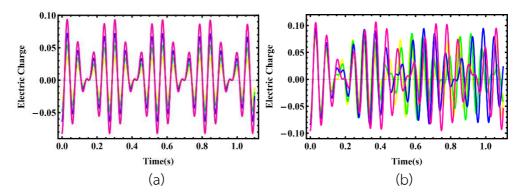


Figure 5: Illustration showing the relation parameter between charges depends on time where δ and σ have slightly different value $(\sigma > \delta)$. (a) Representation the relation between the time-dependent electric charge and time when the initial applied voltage force (E_0) value is unable, (b) Representation the relation between the time-dependent electric charge and time when the resistance (R_{res}) value is unable.

DISCUSSION

From figure 2 (a), we set $q_0=0.5\,\mu C$, $\eta=0.1$, $L_{ind}=450\,H$, $E_0=350Vol$, $R_{res}=250\,W$ are the control variable. But the capacitor is an independent variable and the time-dependent electric charge is a dependent variable. If the capacitance is higher, the wavelength of the vibrating electric charge is longer, but the frequency of the vibrating electric charge is lower. From figure 2 (b), we set $q_0=0.5\,\mu C$, $\eta=0.1$, $C_{Cap}=8\mu F$, $E_0=350Vol$, $R_{res}=250W$ is control variable. The inductance is an independent variable and the time-dependent electric charge is a dependent variable. If the inductance is higher, the amplitude of the vibrating electric charge is higher, and

the wavelength of it is longer. From Figure 3 (a), we set $q_0 = 0.5 \,\mu C$, $\eta = 0.1$, $L_{ind} = 40 \,H$, $C_{\it Cap2} = 8 \mu F$, $R_{\it res} = 100 \, {
m W}$ as the control variable. The initial voltage force $E_{\it 0}$ is the control variable and the time-dependent electric charge is the dependent variable. If E_0 higher, the vibrating of electric charge is smaller, more of the amplitude of the vibrating electric charge is split. From figure 3 (b), we set $q_0 = 0.5\,\mu C$, $\eta = 0.1$, $L_{\scriptscriptstyle ind}$ = 150 H , $C_{\scriptscriptstyle Cap2}$ = $8\mu F$, $E_{\scriptscriptstyle 0}$ = 250 Vol as the control variable. The resistance $R_{\scriptscriptstyle res}$ is the independent variable and the time-dependent electric charge is the dependent variable. If the resistance is higher, the amplitude of the vibrating electric charge is lower. Thus, we can compare the resistance as the damping coefficient of the electric charge. From figure 4 (a), we set $q_0 = 0.5\,\mu C$, $\eta = 0.1$, $L_{ind} = 1.102\,H$, $E_0 = 220 Vol$, $R_{\mbox{\tiny res}} = 100 \, {
m W}$ is control variable. The capacitance is the independent variable and the time-dependent electric charge is the dependent variable. If δ and σ have slightly different, the time-dependent electric charge behaves itself like wave group. If the capacitance is higher, the amplitude of the vibrating electric charge is higher. From figure 4 (b), we set $q_{\scriptscriptstyle 0} = 0.5\,\mu C$, $\eta = 0.1$, $C_{\scriptscriptstyle Cap} = 420\,\mu F$, $E_{\scriptscriptstyle 0} = 220 Vol$, $R_{\scriptscriptstyle res} = 100\,\mathrm{W}$ as control variable. The inductance is the independent variable and the time-dependent electric charge is the dependent variable. If the inductance is higher, the amplitude of the vibrating electric charge is lower. From figure 5 (a), we set $q_0 = 0.5 \,\mu C$, $\eta = 0.1$, $L_{ind}=1.12\,H$, $C_{Cap2}=420\mu F$, $R_{res}=100\,\mathrm{W}$ as the control variable. The initial voltage force $\boldsymbol{E_{0}}$ is the independent variable and the time-dependent electric charge is the dependent variable. If E_0 is higher, the amplitude of the vibrating electric charge is higher. From figure 5 (b) we set $q_0 = 0.5 \,\mu C$, $\eta = 0.1$, $L_{ind} = 1.12 \,H$, $C_{Cap2} = 420 \,\mu F$, $E_{\scriptscriptstyle 0}$ = 220Vol is the control variable. The resistance $R_{\scriptscriptstyle res}$ is the independent variable and the time-dependent electric charge is the dependent variable. If resistance is higher, the amplitude of the vibrating electric charge is higher.

We had explicated the series of RLC loop circuit of resonant system that can be used to effectually substantiated, in a very visual way, the phase relation between the voltages across reactive and resistive elements (Sokol *et al.*, 2013). We have explicated the series of RLC loop circuit of resonant system that can be used to effectually substantiate new soft ferromagnetic magnetic materials made possible the

use of the magnetic amplifier technology in designing competitive electric-power engineering supplies (Eloisa & Romeo, 2004). We can see the time-dependent electric charge perturb via applied voltage force emf has the behavior like wave group but the time-dependent electric charge in paper of title application of linear differential equation in an analysis transient and steady response for second order RLC closed series circuit of Ahammodullah Hasan has not behavior like wave group, if we add the time, the amplitude slowly decreases. A modeling approach of a magnetic amplifier based on the magnetic hysteresis loop of the core soft ferromagnetic material is presented here for a common amorphous magnetic alloy.

CONCLUSIONS

From equation (27), we known that if the value of the parameter of δ is defined as ratio between resistance with double inductance and the free natural angular frequency σ have more different values, the electric charge have their behavior like underdamped. If the value of the parameter of δ is defined as ratio between resistance with double inductance and the free natural angular frequency σ have slightly different values, the electric charge will have their behavior like wave group.

ACKNOWLEDGMENT

We would like to thank you to the Institute for the promotion of teaching science and technology (IPST) of wittayanukulnaree school, phetchabun rajabhat university, science center laboratory and physics division and faculty of science and technology phetchabun rajabhat university for partial support.

REFERENCES

Ahammodullah, H., Md-Abdul, H. and Md. Al-Amin M. (2019). Application of Linear Differential Equation in an Analysis Transient and Steady Response for Second

- Order RLC Closed Series Circuit. *American Journal of Circuits, Systems and Signal Processing*, *5*(1), 1-8.
- Atamp, A. (1990). *Introduction to classical mechanics*. Prentice-Hall, United States America.
- Dino, K. (2019). Negative capacitance or inductive loop-A general assessment of a common low frequency impedance feature, *Electrochemistry Communications*, 98, 58-62.
- Eloisa, G.C. and Romeo, O. (2004). A new passivity property of linear RLC circuits with application to Power Shaping Stabilization. In *Proceedings of the American Control Conference*, pp.1428-1433.
- Goldstein, P. and Safko, J. (2002). *Classical mechanics*. 3th ed., Prentice-Hall, United States America.
- Hutem A. and Masoongnoen N. (2021), SOLVE THE SERIES OF RLC CIRCUIT VIA THE TIME-DEPENDENT VOLTAGE SINE FUNCTION USING WRONSKIAN'S OF DIFFERENTIAL EQUATION. *PSRU Journal of Science and Tehnology*, 6(2), 22-35.
- Kishore L.K. (2008). *Electronic circuit analysis*. second edition, BS Publication.
- Mohazzab, J., Hussain, A., Muhammad, Q. and Mohsin, S. (2008). RLC Circuit Response and Analysis (Using State Space Method). *IJCSNS International of Computer Science and Network Security, 8*(4), 48-54.
- Riley, K.F. and Hobson, M.P. (2006). *Mathematical Methods for Physics and Engineering*. (3th ed.). New York: Cambridge University Press.
- Sadri, H. (1991), Foundations of mathematical physics. Prentice-Hall, United States America
- Sokol P.E., Warren G., Zheng, B. and Smith, P. (2013). A circuit to demonstrate phase relationships in RLC circuits. *Physics Education*, *48*(3), 312-318.
- Sonam, P. (2015). Analyzing the Response of an RLC Circuit. *International Journal of Scientific and Research Publications*, *5*(7), 1-2.
- Susan, M. (2004). *Mathematics for physicsts*. Thomson Brooks, United States America.

- Teoh A.K. and Rahifa, R. (2018). COMPARISON OF NUMERICAL TECHNIQUES IN SOLVING TRANSIENT ANALYSIS OF ELECTRICAL CIRCUITS. *ARPN Journal of Engineering and Applied Sciences*, *13*,(1), 314-320.
- Tikjha, W., Normai, T., Jittburus, U., and Pumila, A. (2018). Periodic with period 4 of a piecewise linear system of differential equations with initial conditions being some points on positive y axis. *PSRU Journal of Science and Tehnology, 3*(2), 26-34.